

Threefold Way for Typical Entanglement (arXiv:2410.11309)



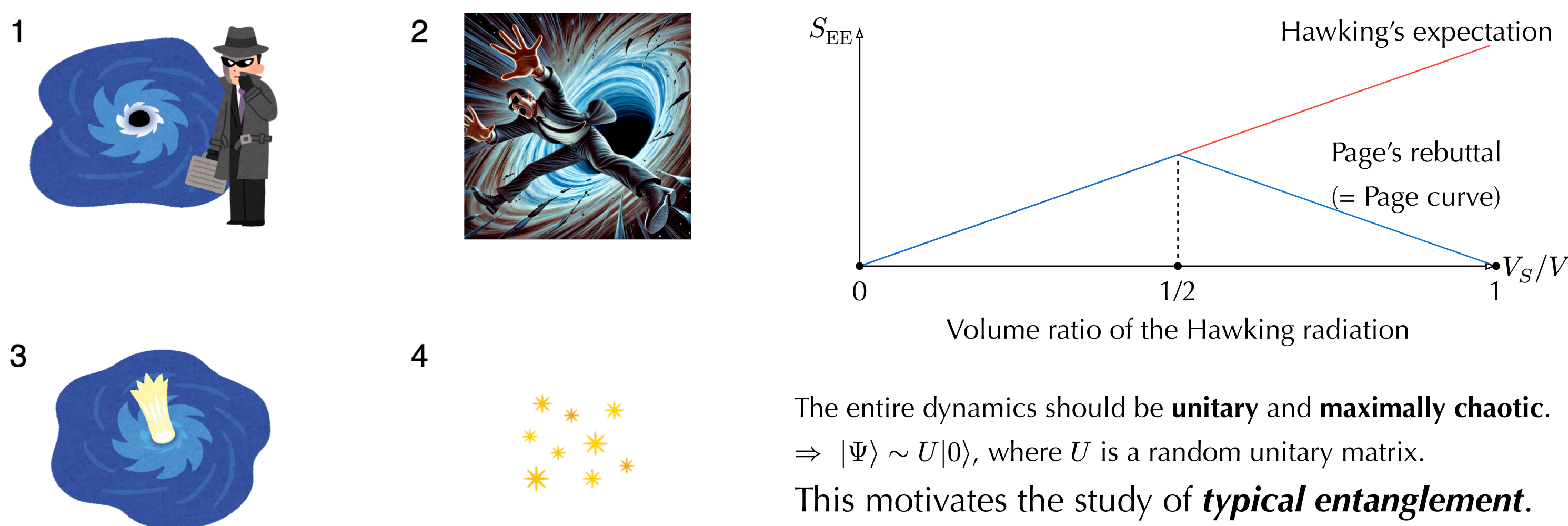
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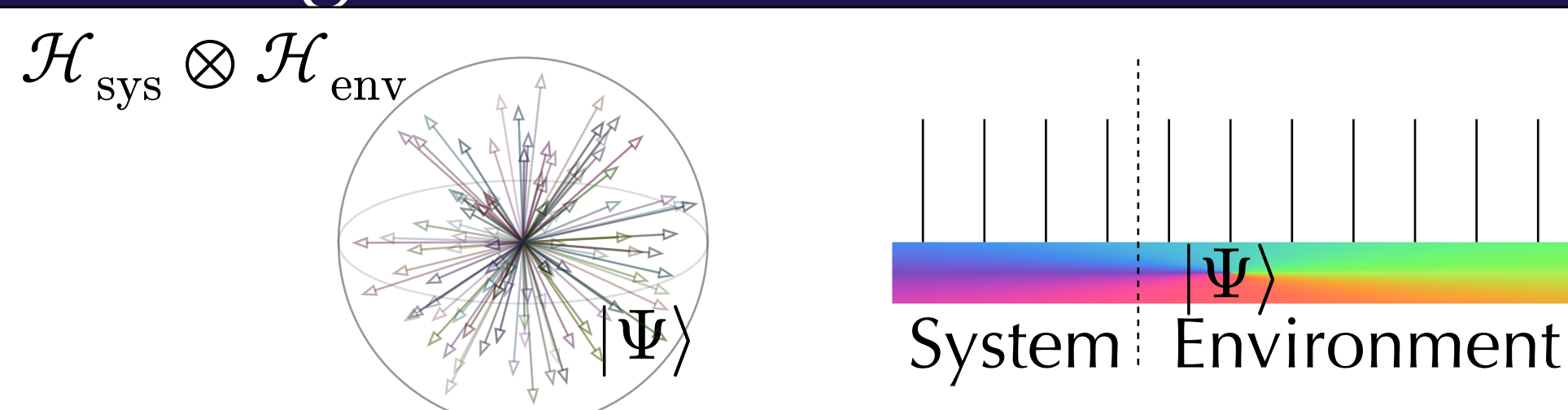
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1. The Black Hole Information Paradox

Thermalization is inconsistent (?) with unitary dynamics.



2. Typical Entanglement



- **Typical** \approx Haar-random (uniformly random sampling on the entire Hilbert space)
- **Entanglement spectrum** = eigvals of the reduced density matrix
- Random matrix theory is useful to evaluate typical entanglement[1–4].

3. Laguerre Unitary Ensemble of Random Matrices

Preparing Haar-random state from the entire Hilbert space:

$$|\Psi\rangle \sim \text{Haar} \Leftrightarrow \rho_{\text{sys}} = \mathbf{W} \mathbf{W}^\dagger.$$

where i.i.d. $W_{\text{sys,env}} \sim \mathcal{CN}(0, 1)$, $|\Psi\rangle = \sum_{\text{sys,env}} W_{\text{sys,env}} |\text{sys,env}\rangle$ [5,6].

$\Rightarrow \rho_{\text{sys}}$ follows the **Laguerre unitary ensemble (LUE)** of RMT.

The joint probability p distribution of eigenvalues $\{\lambda_i\}$ of ρ_{sys} is **Laguerre distribution**:

$$p(\lambda_1, \dots, \lambda_m) \propto \prod_{i=1}^m \lambda_i^{\frac{\beta}{2}(n-m+1)-1} e^{-\frac{\beta}{2}\lambda_i} \prod_{1 \leq i < j \leq m} |\lambda_i - \lambda_j|^\beta$$

with $\beta = 2$ for LUE.

4. TRS and Threefold Way

What if we change complex random variables to real random variables?

$\Rightarrow |\Psi\rangle$ and ρ_{sys} obtain **time reversal symmetry (TRS)** of $\mathcal{T} = K$:
 $|\Psi\rangle \rightarrow \mathcal{T}|\Psi\rangle = K|\Psi\rangle = |\Psi\rangle,$

$$\rho_{\text{sys}} \rightarrow \mathcal{T} \rho_{\text{sys}} \mathcal{T}^{-1} = K \rho_{\text{sys}} K = \overline{\rho_{\text{sys}}} = \rho_{\text{sys}}.$$

Note that there is the other & nonequivalent kind of TRS!

Time Reversal Symmetry

$$\begin{aligned} \text{Integer spin: } \mathcal{T}_+^2 &= +\mathbb{1}. & \text{ex. } \mathcal{T}_+ &= K \\ \text{Half-integer spin: } \mathcal{T}_-^2 &= -\mathbb{1}. & \text{ex. } \mathcal{T}_- &= \sigma_y K \end{aligned}$$

Imposing TRS: $\mathcal{T} \rho_{\text{sys}} \mathcal{T}^{-1} = \rho_{\text{sys}}$ allows threefold way of Laguerre ensemble:

- $\beta = 1$ (Laguerre orthogonal ensemble)
- $\beta = 4$ (Laguerre symplectic ensemble).

5. Prohibition of $\mathcal{T}_-^2 = -\mathbb{1}$ TRS Eigenstate

Kramers' theorem: $\mathcal{T}_-^2 = -\mathbb{1}$ TRS cannot have the eigenstate $\mathcal{T}_-|\Psi\rangle = |\Psi\rangle$.

Proof: \mathcal{T} is anti-unitary, thus $\langle \mathcal{T}_- a | \mathcal{T}_- b \rangle = \overline{\langle a | b \rangle}$.

$$\langle \psi | \mathcal{T}_- \psi \rangle = \langle \mathcal{T}_- \psi | \mathcal{T}_-^2 \psi \rangle = \langle \mathcal{T}_-^2 \psi | \mathcal{T}_- \psi \rangle = -\langle \psi | \mathcal{T}_- \psi \rangle$$

This implies $\langle \psi | \mathcal{T}_- \psi \rangle = 0$, thus $|\psi\rangle$ is orthogonal to $\mathcal{T}_-|\psi\rangle$. \square

6. Key Questions

Threefold way of Laguerre Ensemble[7]

	$\mathcal{T}_+^2 = +\mathbb{1}$ ($\mathcal{T}_+ \approx K$) $\beta = 1$	No TRS $\beta = 2$	$\mathcal{T}_-^2 = -\mathbb{1}$ ($\mathcal{T}_- \approx \sigma_y K$) $\beta = 4$
Laguerre	orthogonal,	unitary,	symplectic ensemble
	LOE	LUE	LSE
$\mathcal{T} \Psi\rangle = \Psi\rangle$	real vector	complex vector	<i>unknown</i>

1. Possible to construct the pure state that shows $\rho_{\text{sys}} \sim \text{LSE}$?
2. Beyond threefold way if general symmetries?

7. Key Idea: Symmetry Fractionalization

Fractionalization of TRS

$$\mathcal{T}^2 = [\mathcal{T}_{\text{sys}} \otimes \mathcal{T}_{\text{env}}]^2 = +\mathbb{1}$$

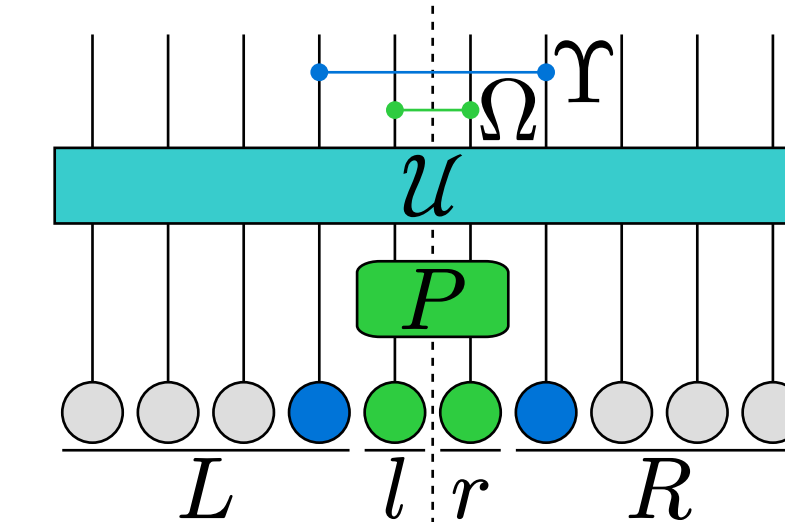
$$\swarrow \mathcal{T} = \Upsilon (\mathcal{T}_{\text{sys}} \otimes \mathcal{T}_{\text{env}}) \Upsilon^\dagger \searrow$$

$$\mathcal{T}_{\text{sys}}^2 = -\mathbb{1}$$

$$\mathcal{T}_{\text{env}}^2 = -\mathbb{1}$$

- We found $\Upsilon = \frac{1-i}{2} [\mathbb{1}_4 - i(\sigma_y \otimes \sigma_y)]$ and proved $\Upsilon|\Psi\rangle$ follows LSE.
- Similarly, $\Omega = \sum_{g_l, g_r} \omega(g_r, g_r^{-1} g_l) |g_l, g_r\rangle \langle g_l, g_r|$ fractionalizes general symmetries: $\Omega[D(g) \otimes D(g)]\Omega^\dagger = D(g) \otimes \overline{D(g)}$, $D(g)D(g') = \omega(g, g')D(gg')$.

8. General Setup



- $G = G_0$ or $G_0 \rtimes \mathbb{Z}_2^{\mathcal{T}}$ (G_0 : unitary, $\mathbb{Z}_2^{\mathcal{T}}$: anti-unitary).
- l, r : $|G_0|$ -dimensional qudit (regular representation of G_0 : $\langle g | g' \rangle = \delta_{g, g'}$).
- P : Projection of $|G_0|^2$ -dimensional Hilbert space of $l \cup r$ onto G_0 -symmetric $|G_0|$ -dimensional basis $\forall g \in G_0$, $|\psi_g\rangle = \frac{1}{\sqrt{|G_0|}} \sum_{h \in G_0} |hg\rangle |h\rangle$.
- $\mathcal{U} \sim$ Haar measure on the projected $d_L d_R |G_0|$ -dimensional space:
 $|\Psi\rangle = \sum_{L, g \in G_0, R} c_{L, g, R} |L\rangle |\psi_g\rangle |R\rangle = \frac{1}{\sqrt{|G_0|}} \sum_{L, g_l, g_r, R} c_{L, g_r^{-1} g_l, R} |L\rangle |g_l\rangle |g_r\rangle |R\rangle$
- Ω fractionalizes G_0 , Υ fractionalizes $\mathbb{Z}_2^{\mathcal{T}}$.

9. Results and Conclusions

Entanglement-spectrum statistics of...

$$\begin{aligned} G = G_0 \rtimes \mathbb{Z}_2^{\mathcal{T}} \text{ is } & \left[\bigoplus_{\alpha: R_1} \frac{\mathbb{1}_{d_\alpha}}{d_\alpha} \otimes \text{LOE}_{\alpha}^{d_L d_\alpha \times d_R d_\alpha} \right] \\ G = G_0 \text{ is } & \left[\bigoplus_{\alpha} \left[\frac{\mathbb{1}_{d_\alpha}}{d_\alpha} \otimes \text{LUE}_{\alpha}^{d_L d_\alpha \times d_R d_\alpha} \right] \right] \\ & \oplus \left[\bigoplus_{\alpha: R_0} \frac{\mathbb{1}_{2d_\alpha}}{d_\alpha} \otimes \text{LUE}_{\alpha}^{d_L d_\alpha \times d_R d_\alpha} \right] \\ & \oplus \left[\bigoplus_{\alpha: R_{-1}} \frac{\mathbb{1}_{d_\alpha}}{d_\alpha} \otimes \text{LSE}_{\alpha}^{d_L d_\alpha \times d_R d_\alpha} \right]. \end{aligned}$$

Until our work

- The setup which follows LSE have been elusive.
- What this work revealed are:
- The LSE setup can be constructed by fractionalizing TRS of the LOE setup.
- Extended the setup to general symmetries.
- Entanglement-spectrum statistics is **direct sum of the threefold way**.

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